

# Characteristics of rock breaking with particle jet impact in deep geothermal reservoir

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## Abstract

Particle jet impact can fracture the rock, which can be utilized to improve the rock breaking efficiency of deep hot dry rock. In order to further explore the rock breaking characteristics, the breaking mechanism of hot dry rock under the action of particle jet impact were analyzed based on the Smoothed Particle Hydrodynamics (SPH) method, considering the deep geothermal high temperature and high confining pressure. The results show that the tensile stress generated by particle impact is much larger than the tensile strength of rock, and the rebound effect constituted an important factor in the crater diameter determination. Both the impact depth and breaking volume increased with the increase in the jet inlet velocity. Both the impact depth and breaking volume increases with the increase in the diameter. When the particle concentration exceeded a critical value (3%), both the impact depth and breaking volume increased slowly. The confining pressure will hinder the efficiency of particle impact rock breaking. High temperature will generate larger impact depth and breaking volume. The research can provide theoretical supports and parameter guidance for the application of efficient deep hot dry rock breaking technology with particle impact.

**Keywords:** Deep formation; Hot dry rock; Particle jet impact; Rock breaking efficiency; Rock breaking mechanism

## 1. Introduction

The particle jet is utilized in rock cutting [1,2] or in the shot blasting treatment technology [3,4] and is similar to the abrasive waterjet. However, spherical steel particles with a greater diameter (ranging from 0.5 to 5 mm) and lower particle concentration (ranging from 0.5 to 5% by volume) are utilized in the particle jet, which make it different from abrasive waterjet [5,6]. The steel particles act as an important factor in the erosion of rocks with a waterjet. The particles erupting from the nozzle at a high speed and frequency impact the rock within an extremely short time period and upon an extremely low sized contact area. The instantaneous contact stress is significantly high, thus changing the conventional rock breaking trend and significantly increasing the energy utilization [7,8].

The particle jet cutting performance is a nonlinear impact dynamics problem affected by various factors such as jet pressure, standoff distance, particle concentration, and the material. Moreover, the target material can be removed from the particle impact accelerated by water at a high velocity [9–11]. The solid particle erosion consists of the following two parts: deformation damage removal and cutting removal [12]. The rock is penetrated by particles with a high instantaneous dynamic load, high deformation, and short dwell time [13]. By the utilization of both the Newton motion equation and law of Meyer, the equation of the white iron crater volume penetration by particles was established, whereas the abrasive cut depth was proportional to the abrasive kinetic energy [14]. Experiments were also carried out for analyzing the effects of the cavity height on cavity erosion under a gas–solid flow. The spatial distribution of particles within the cavity and erosion related to particle

impacts for the variable cavity aft wall heights and for two particle sizes (44 and 106  $\mu\text{m}$ ) were studied. The maximum erosion rate was quite significant for higher sized particles [15]. A comprehensive modeling and computational study was presented for the erosion in elbows determination due to sand particles being entrained in air both by simulation and experimental methods, indicating that as the velocity increased, the metal loss increased following a power law [16].

With the objective of identifying and revealing the practical applications of a particle jet, investigation of the effects of the waterjet parameters on rock breaking performance is extremely important and highly desirable. The cutting depth efficiency tends to increase with increasing water pressure and traverse speed and with decreasing standoff distance. The cutting volume efficiency is strongly dependent on the standoff distance [17]. According to literature, with the contribution of the abrasive water jet, the drilling depth increased by  $\sim 63\%$ , the thrust force and torque were reduced by  $\sim 15$  and  $20\%$ , respectively, whereas the bit wear reduced significantly [18]. Both the theoretical analysis and experimental measurements demonstrated that the volume of the crater created by individual particle impingement, with the impingement velocity having been ranged between 5 and 11 m/s, was linearly proportional to the particle kinetic energy [19]. The kinetic energy of the wear particles that were removed increased linearly along with the erosion depth [20]. As the extent of strains exceeded the elastic limit, the impacted solid particles skidded over the surface. The threshold strain was theoretically derived as a function of the impact velocity and the rotating velocities as well as the particle surface contact duration [21]. A phenomenological model of the three-phased flow inside an abrasive waterjet machining cutting head was presented. The model was validated by an extensive set of experimental data with wide variations in cutting-head geometry, operating pressure, and abrasive mass flow rates [22]. An analytical model considering the size and shape of the abrasive particles was developed for the prediction of the brittle polycrystalline material total cut depth [23]. The tensile fracture caused by the shock pressure wave of abrasion was the main reason for the breaking of the rock [24,25]. A framework for sand collision modeling was presented including the effects of elastic, elastic-plastic and plastic deformations, surface adhesion, and size dependent property variations of sand [26]. The effects of pressure confinement on rock damage were analyzed by the ANSYS implicit-explicit solution [27]. The rock breaking caused by particle impact was investigated by the one-dimensional stress wave theory [28]. The abrasive wear is a rapid and severe process due to the contact between the abrasives and the solid material surfaces. This type of wear is usually classified into the following two categories: (i) the 2-body abrasive wear and (ii) the 3-body abrasive wear [29]. An integral approach with consideration of both the plastic and elastic-plastic material removal modes was presented. The erosion was characterized by two failure modes, the intercrystalline fracture and lateral cracking, whereas the power exponent regarding the particle velocity effects was 2.44 [30,31].

The smoothed particle hydrodynamics (SPH) simulations were utilized for particle embedment simulation. Regarding a given particle size, the embedment increased along with the incident velocity, whereas a critical minimum velocity, below which embedment would not occur, was predicted [32]. An SPH simulation of a cone penetrometer test in soft, cohesive soil was implemented. The higher soil strengths could be reproduced in the simulation if an improved shear model including the non-negligible friction angle was provided [33]. The energy absorption behavior of empty and aluminum metal foam-filled tubes with various taper angles was evaluated by the Finite Element Method (FEM) coupled with SPH [34]. A coupled FEM/SPH modeling technique was investigated for the numerical simulation of the quasi-static axial crushing of polystyrene foam-filled aluminum thin-walled aluminum tubes [35]. An FE methodology was developed for tensile stresses in the coating surface under a single particle impact analysis simulating the particulate erosion conditions [36].

Besides, most previous studies focused on rock cutting with the pure waterjet or abrasive waterjet, and not the particle jet, being different in utilization of spherical steel particles with a higher diameter and lower particle concentration. In a typical abrasive waterjet, the shape of the abrasives is irregular, and a garnet or quartz is usually utilized as the abrasive with diameter ranging from 0.06 to 0.3 mm. These characteristics of the abrasives

lead to the serious wearing of the parts (such as nozzles and pipes) of an abrasive waterjet. In contrast, during a particle jet, the equipment wears only slightly. In most previous studies, the effects of waterjet were not considered, and the FEM/SPH method was rarely utilized for the investigation of the particle jet rock breaking problems. This technology can significantly enhance the hot dry rock breaking efficiency owing to the higher kinetic energy of particles. In practice, a multi large spherical steel particle impact is required for the rapid cutting of the rock. Furthermore, in most previous studies only the theoretical calculation or the experimental method was utilized for the investigation of the rock breaking performance, and the effects of high confining pressure and temperature were seldom studied.

In this study, a numerical simulation was utilized for the establishment of the relationship between the rock breaking performance and particle parameters for hot dry rock, with consideration of the effects of waterjet based on the SPH method, conforming to practical operational conditions in an effort for providing an efficient design for practical application.

## 2. Finite element model

### 2.1. Smoothed Particle Hydrodynamics-Finite Element Method

In the SPH method, the continuous fluid is dispersed to interact with particles, and the change in particles is utilized to describe the change in the continuous fluid by the Diffuse Element Method. The SPH method is widely utilized in the high deformation of penetration study [32–35]. In this study, the SPH coupled with the FEM was applied to investigate the relationship between the particle jet parameters and the rock breaking performance (depth and volume). Furthermore, both the SPH and FE methods were utilized for the particle jet and rock models establishment, respectively.

Both the particle and water were modeled by the SPH method, which was a pure Lagrange method and did not require a mesh based upon the interpolation theory. A series of uniformly distributed smoothed particles was utilized having a variety of physical properties for the partial differential equations solution under various conditions. The kernel estimated value  $f(x)$  could be expressed by using Eq. (1) [36–39] as follows:

$$f(x) = \int_{\Omega} f(x') W(x - x', h) dx' \quad (1)$$

where  $f(x)$  is the function of the 3D coordinate vector  $x$ ,  $\Omega$  is the support domain of the point  $x$ ,  $x - x'$  is the distance between the particles,  $h$  is the smooth length of the SPH particle changing along with time and space, and  $W(x - x', h)$  is the kernel function, usually defined by the auxiliary function ( $\theta(x - x')$ ):

$$W(x - x', h) = \frac{1}{h(x - x')^d} \theta(x - x') \quad (2)$$

where  $d$  is the space dimensionality.

The cubic b spline curve is most commonly utilized for a smooth function in the SPH method as described by the following equation:

$$\theta(u) = C_b * \begin{cases} 1 - \frac{3}{2}u^2 + \frac{3}{4}u^3, |u| \leq 1 \\ \frac{1}{4}(2 - u)^3, 1 \leq u \leq 2 \\ 0, |u| \geq 2 \end{cases} \quad (3)$$

where  $C_b$  is the normalization constant, confirmed by the space dimension  $u = \frac{x - x'}{h}$ .

Substitution of Eq. (2) into Eq. (1) and the conversion of the continuous form integral equation into the discrete form equation led to the following equation:

$$f(x) = \sum_{i=1}^n \frac{m_i}{\rho_i} \frac{1}{h(x - x'_i)^d} \theta(x - x'_i) \quad (4)$$

where  $\rho_i$  is the density of particle  $i$  and  $m_i$  is the mass of particle  $i$ .

## 2.2. Water and Rock Damage Model

The 1# elastic material model in the ANSYS-DYNA was utilized for particles, and the parameters are listed in Table 1 [37,38].

**Table 1**

1# elastic material model parameters for particles [37,38]

Material	$\rho_p$ [g/cm <sup>3</sup> ]	$\nu_p$	$E_p$ [GPa]
Particle	7800	0.3	210

The MAT\_NULL material model was utilized for water. The GRUNEISEN state equation expressed by Eq. (5) was also utilized for water and the parameters are also listed in Table 2 [40–43]:

$$P_{MG} = \frac{\rho_0 C_{MG}^2 \left[ 1 + \left( 1 - \frac{\gamma_0}{2} \right) \mu - \frac{a}{2} \mu^2 \right]}{\left[ 1 - (S_1 - 1) \mu - S_2 \frac{\mu^2}{\mu + 1} - S_3 \frac{\mu^3}{(\mu + 1)^2} \right]} + (\gamma_0 + a \mu) E a \quad (5)$$

where  $P_{MG}$  is the pressure,  $E$  is the internal energy per unit volume,  $C_{MG}$  is the intercept of the  $u_s - u_p$  curve,  $\rho_0$  is the density,  $S_1$ ,  $S_2$ , and  $S_3$  are the slope coefficients of the curve,  $\gamma_0$  is the GRUNEISEN constant, and  $a$  is the volume correction.

**Table 2**

Water parameters for simulation [43]

Material	$\rho_0$ [g/cm <sup>3</sup> ]	$C$ [m/s]	$S_1$	$S_2$	$S_3$	$a$	$\gamma_0$
Water	1	1480	2.56	-1.986	0.2286	1.397	0.49

The Holmquist Johnson Cook (HJC) model was utilized for the rock, which is usually employed as the rock damage model under a high deformation, high strain rate, and high pressure, with consideration of the effects of rock material damage, strain rate, and hydrostatic pressure to the yield force [44,45]. The yield surface equation of the HJC model is as follows:

$$\sigma^* = [A(1 - D) + BP^{*N}](1 + C \ln \bar{\varepsilon}^*) \quad (6)$$

where  $\sigma^*$  is the equivalent stress which is dimensionless and obtained from the actual equivalent stress divided by the static compressive strength  $f'_c$ ,  $P^*$  is the hydrostatic pressure which is dimensionless and obtained from the actual hydrostatic pressure divided by the static compressive strength  $f'_c$ , and  $\bar{\varepsilon}^*$  is the strain rate that is

dimensionless and obtained from the actual strain rate divided by the reference strain rate  $\dot{\varepsilon}_0$ .  $D$  is the damage ratio and  $A$ ,  $B$ ,  $N$ , and  $C$  are the material strength parameters.

The damage was described by the accumulation of the equivalent plastic and plastic volumetric strains. The damage evolution equation is represented by Eq. (7) [46,47]:

$$D = \sum \frac{\Delta \varepsilon_p + \Delta \mu_p}{\varepsilon_p^f + \mu_p^f} \quad (7)$$

where  $\Delta \varepsilon_p$  and  $\Delta \mu_p$  are the equivalent plastic and plastic volumetric strains, respectively, in an iteration, and

$\varepsilon_p^f$  and  $\mu_p^f$  are the equivalent plastic strain and plastic volumetric strain, respectively, under the atmospheric pressure. The parameters of the HJC model for the rock during the simulation are listed in Table 3 [46].

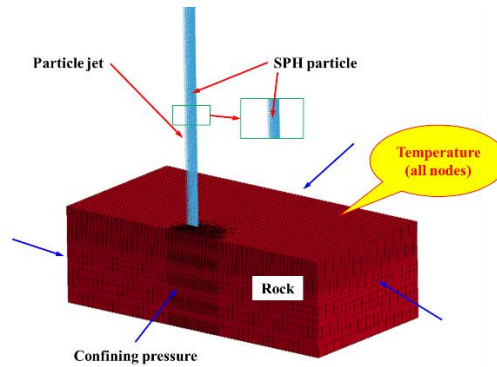
**Table 3**

HJC model parameters for the rock during simulation [46]

$\rho$ [kg/m <sup>3</sup> ]	$G$ [GPa]	$f'_C$ [MPa]	$A$	$B$	$C$	$N$	$S_{MAX}$	$D_1$	$D_2$
2600	20.8	61	0.93	1.60	0.08	0.79	7.0	0.04	1.0
$\varepsilon_{fmin}$	$T$ [MPa]	$p_c$ [MPa]	$\mu_c$	$p_1$ [GPa]	$\mu_1$	$k_1$ [MPa]	$k_2$ [MPa]	$k_3$ [MPa]	$\dot{\varepsilon}_0$
0.004	4	2.03	0.66	1.0	0.072	85	-171	208	1

### 2.3. Finite element model

The 3D FE model established for a particle jet is presented in Fig. 1. The grid number of the rock utilized in this case was approximately 180,000. The SOLID 164 element with 8 nodes was applied for rock meshing. The number of SPH particles for the steel particles was proportional to the SPH particles for water that could be determined by the designed particle concentration (the particles concentration is 1% by volume in Fig. 1). The NON-REFLECTING boundary condition was utilized to constrain the underside and the four side faces of the rock. The contact type to the rock was set as \*CONTACT\_ERODING\_NODES\_TO\_SURFACE. The inlet velocities of both water and particles were set as the initial velocities. In particular, the red color in the waterjet was utilized to represent the spherical steel particles as shown in Fig.1. The rock erosion was resulted by the removal of the rock elements. When the equivalent plastic strain exceeded the  $\varepsilon_{fmin} = 0.004$  (the minimum equivalent plastic strain that caused the rock to become broken as listed in Table 3), the FE of the rock began to be removed.



**Fig. 1.** 3D Finite element model for the particle jet based on the SPH-FEM method.

Confining pressure applying steps are as follows: Define the load curve by keyword *\*Define\_Curve*; Set the "Segment" set on the plane where confining pressure is required; Load curve is applied to the "Segment" set by keyword *\*Load\_Segment\_Set*; Define *\*Interface\_Springback\_Lsdyna* keyword. Select all the nodes and apply the temperature loads to hot dry rock by keyword *\*Initial\_Temperature\_node*, and solve and output *\*K* file.

#### 2.4. Particle jet parameters

The rock breaking simulation were performed under various particle jet parameters (particle velocity, particle diameter, concentration, confining pressure, and temperature). The particle velocity was controlled from 100 to 240m/s. The particle diameter ranged from 0.5 to 2.5 mm, whereas the particle concentration within the waterjet was changed from 0.5 to 5% by volume. The confining pressure ranged from 0 to 50 MPa. The temperature for the hot dry rock varied from 100 to 400°C. The simulation parameters for the particle jet impact are listed in Table 4.

**Table 4**

Particle jet parameters

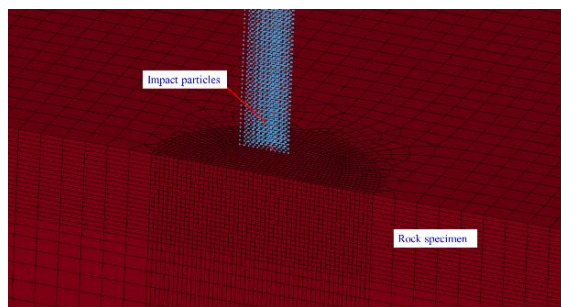
Water pressure ( $P_w$ ) [m/s]	100	136	165	195	220	240
Particle diameter ( $d_p$ ) [mm]	0.5	1	1.5	2	2.5	
Particle concentration by volume ( $C_p$ ) [%]	0.5	1	2	3	4	5
Confining pressure ( $P_c$ ) [MPa]	0	10	20	30	40	50
Temperature ( $T_i$ ) [°C]	25	100	200	300	400	

### 3. Result and discussion

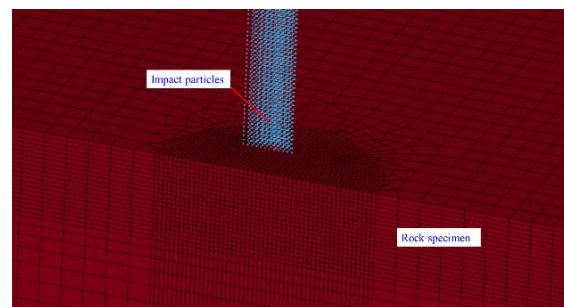
The rock breaking mechanism and effects particle jet parameters on the rock breaking depth and volume were analyzed by numerical simulation.

#### 3.1. Analysis of rock breaking process

The rock breaking by particle jet impact is presented in Fig. 2. The particles with a high velocity were accelerated by the fluid impact upon the rock surface. At a higher erosion depth, the volume of the crater was higher along with the continuous impact of particles. Particles following impact constantly rebounded and left the crater area as presented in Fig. 2. The diameter of the initial formed crater was almost the same as the jet diameter as presented in Fig. 2b. Because the particle diameter was higher, the particle density was higher than the density of ordinary abrasives and the rock was hard to be penetrated, the particles impacted the hard rock and rebounded. Moreover, as well-known, the secondary flow of the fluid carrying particles outside the primary impact zone can intensify this rebound effect. When the dwell time increased, the rebound particles continuously stroke the side face of the craters and the diameter of the craters increased due to the rebound effect as presented in Figs. 2c and d. Therefore, the rebound effect acted as an important factor in the crater diameter determination during the particle jet impact.

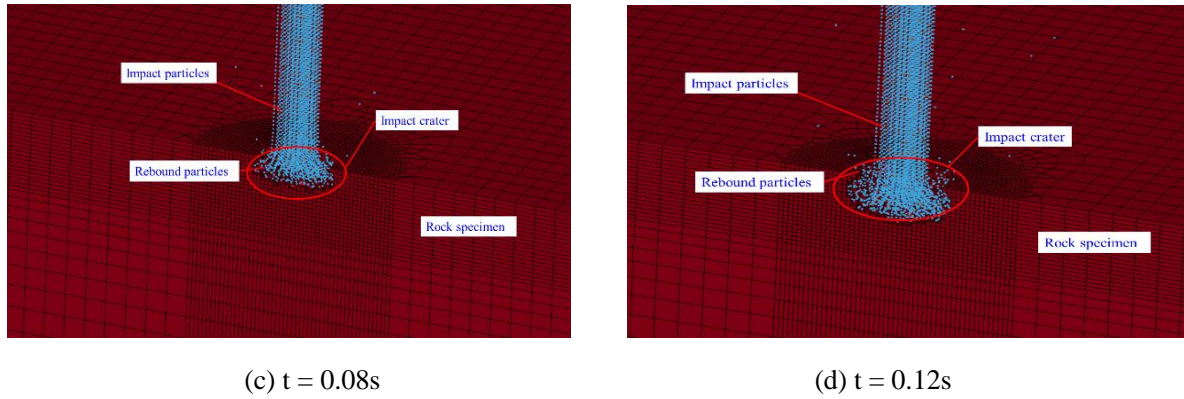


(a)  $t = 0$  s



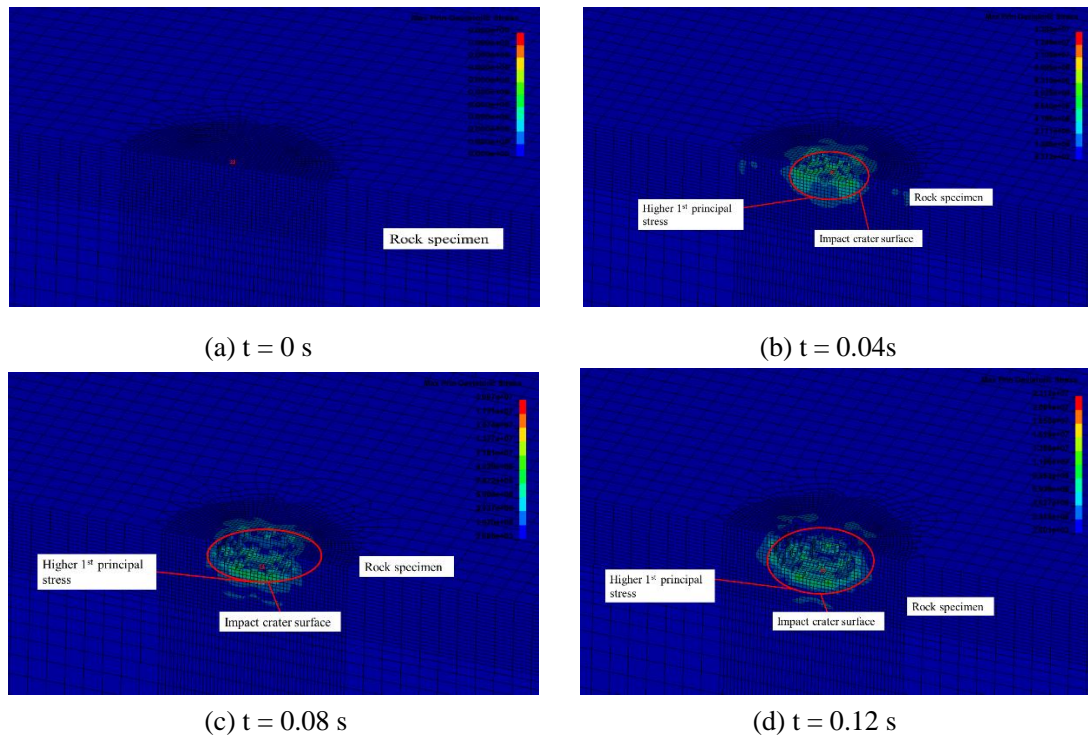
(b)  $t = 0.04$  s





**Fig. 2.** Rock breaking process by particle jet

For brittle materials such as concrete, rocks, and cast iron, the First Strength Theory (or the Tensile Strength Theory) is usually utilized to estimate whether these materials are broken or not. The value of the first strength theory was also the 1<sup>st</sup> principal stress (maximum principal stress) in the ANSYS software. During impact, the 1<sup>st</sup> principal stress distribution of the rock is presented in Fig. 3. The rock began to break when the 1<sup>st</sup> principal stress exceeded the tensile strength of the rock. The maximum value of 1<sup>st</sup> principal stress was 23.1 MPa, significantly higher than the tensile strength of the rock (the tensile strength of marble was  $\sim 4$  MPa) in the time range of 0.04 s; therefore, the rock could be broken easily (Fig. 3b). The higher values of the 1<sup>st</sup> principal stress were distributed on the surface areas of the crater and the 1<sup>st</sup> principal stress became lower when the distance between the points of the rock and the surface of the crater increased as presented in Figs. 3 b–d. The bottom of the crater, where the 1<sup>st</sup> principal stress was higher, was the main broken area of the rock. The 1<sup>st</sup> principal stress was higher only in the side of the crater where the rebounding particles impacted, as presented in Figs. 3c and d. The side of the rock could not be broken because of this rebound impact and this result was well consistent with the results presented in Fig. 2.

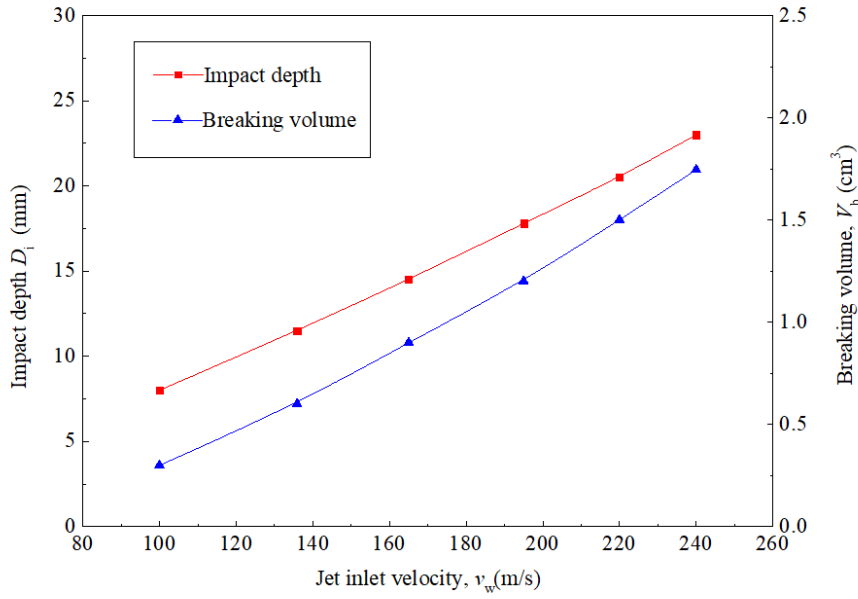


**Fig. 3.** The 1st principal stress distribution.

### 3.2. Effects of particle jet parameters

#### 3.2.1. Jet inlet velocity

The particle jet energy is directly related to the jet inlet velocity ( $v_w$ ). Therefore, the impact depth and breaking volume of the rock increased as the jet inlet velocity increased in agreement with previous findings [9,48]. Fig. 4 demonstrates the effects of jet inlet velocity on the impact depth and volume. The results clearly indicated that the impact depth increased as the  $v_w$  increased. The impact depth and velocity had an approximately linear relationship. The depth of cut increased rapidly with the increase in the kinetic energy of the impacting particles due to the increase in the water inlet velocity [48,49]. As the velocity increased, the breaking volume increased. When the velocity increased, both the kinetic energy of the particles and water would increase significantly, leading to a higher impact energy that acted on the rock, consequently the rock could easily be broken. Therefore, the rock breaking volume increased at a higher water inlet velocity. Besides, due to the higher density of particles and harder rock, it was believed that the energy increase of particles was the main factor that influenced the rock breaking volume changes.

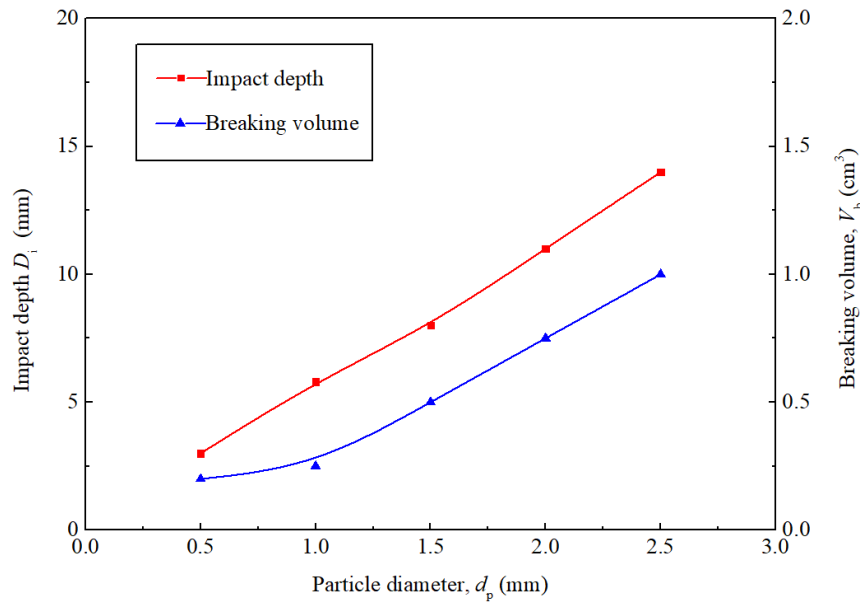


**Fig. 4.** Velocity effect on impact depth and volume.

### 3.2.2. Particle Diameter

The particle diameter ( $d_p$ ) is an important parameter that is directly related to the particle kinetic energy. Particles with a higher diameter can generate a higher energy that impacts the rock. Fig. 5 exhibits that the impact depth increases with the increase in the diameter, and this result is in agreement with previous findings [50–53]. Furthermore the results indicated a nearly linear relationship between the impact depth and particle diameter. This could be explained by the fact that lower sized particles had a lower kinetic energy compared to higher sized particles due to lower masses. Therefore for a single particle, a higher sized particle could cause a higher impact force that would act on the rock due to the higher kinetic energy than lower sized particles. Moreover, lower sized particles were diverged further from the nozzle axis to the exterior and hit on a higher sized region on the rock surface, whereas higher sized particles changed only slightly the corresponding trajectories and hit on a lower sized region encircling the nozzle axis. Consequently, regarding the same particle flow rate, the higher sized particles might cause a high impact concentration (the number of repeated impacts upon the same location) because the impact area was lower sized [52]. The breaking volume increases with the increase in the diameter. When the particle diameter exceeded 1.0 mm, the breaking volume began to increase rapidly due to the higher kinetic energy and higher impact concentration for higher sized particles. Furthermore, the speed increased as the particle diameter increased (the slope of the curve increased as the particle diameter increased). Therefore, the breaking volume of a particle water jet strongly depends on particle diameter.

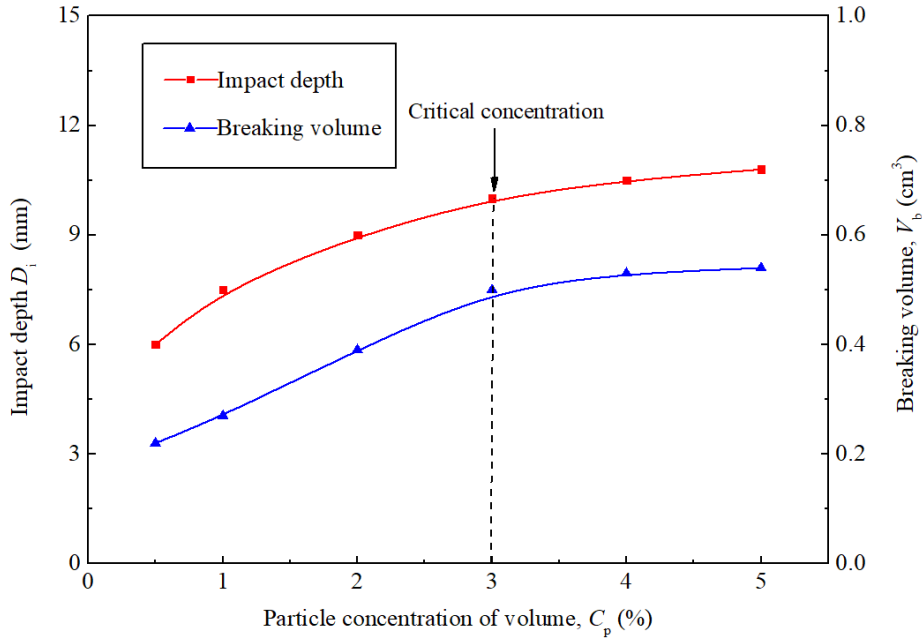




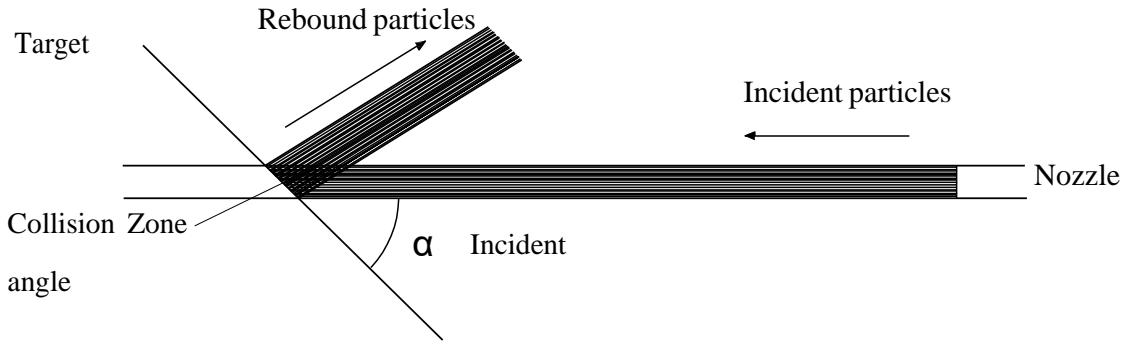
**Fig. 5.** Diameter effect on the impact depth and volume.

### 3.2.3. Particle Concentration

The particle concentration ( $C_p$ ) is defined as the volume concentration of the water flow. The results as presented in Fig. 6 indicate that the impact depth increases as the particle concentration increases. When the particle concentration was below the critical concentration (3%), the impact depth increased rapidly with the increase in the concentration. Besides, when the particle concentration was  $>3\%$ , the impact depth increased slowly. This could be explained by the fact that an interference effect caused by the collisions between the incident and rebound particles existed, which was studied in detail [54–57] as presented in Fig. 7. Furthermore it was quite useful for the critical concentration or mass flow to be estimated for interference effects minimization during rock breaking, in order for a higher proportion of the stream energy leaving the nozzles to reach the intended target. It could be confirmed that the critical concentration was 3% in this study. When the particle concentration was  $<3\%$ , the probability of collisions was low, and most particles arrived at the target surface without sustaining an inter-particle collision, therefore the impact depth increased dramatically. When the particle concentration was  $>3\%$ , the percentage of particles without collision reduced dramatically, possibly resulting in the significant decrease in the particle kinetic energy eroding the target, therefore the impact depth increased slowly. Moreover, as the incident angle was higher ( $\alpha = 90^\circ$ ), the collision zone size was higher, and therefore the probability of survival was quite lower than that at a lower incident angle [55]. The interference effect was quite complicated; therefore, the degree of particle–particle interactions will be investigated in the future. The breaking volume also increases with the increase in the particle concentration, similar to the effects of concentration on the impact depth. When the particle concentration was below the critical concentration (3%), the breaking volume increased rapidly with the increase in the concentration. Besides, when the particle concentration was  $>3\%$ , the breaking volume increased slowly owing to interference caused by the collisions between the incident and rebound particles.



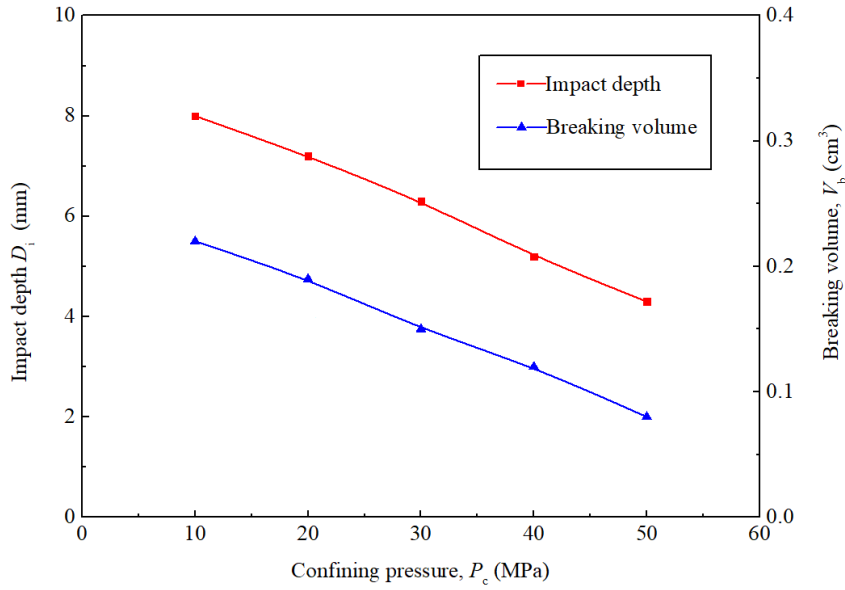
**Fig. 6.** Concentration effect on the impact depth and volume.



**Fig.14.** Schematic diagram of collisions between incident and rebound particles [53].

#### 3.2.4. Confining Pressure

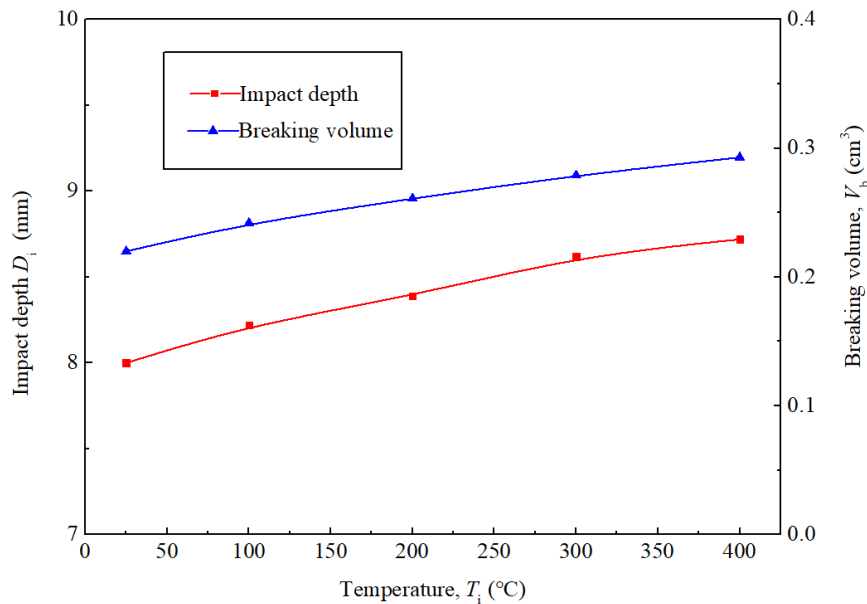
The relationship between the impact depth and breaking volume and the standoff distance for various confining pressure is presented in Fig. 7. With the increase of confining pressure, the impact depth and breaking volume decrease. When the high pressure particle jet contacts the rock, the energy is exchanged rapidly, and the particle impact produces huge impact force on the rock surface. When the strength of tensile and shear strength of the rock is exceeded, the rock is broken. Because the tensile strength of the rock will be improved correspondingly under high confining pressure, the damage development of the rock is slow under high confining pressure. In addition, part of the energy of particle impact will propagate in the rock in the form of stress waves, resulting in tensile effect on the rock. High confining pressure will affect the peak value and frequency of stress waves in the rock, thus weakening the damage effect of stress waves on the rock. The rock damage and breakage time required will increase. Therefore, with the increase of confining pressure, the impact depth and breaking volume decrease.



**Fig. 7.** Confining pressure effect on the impact depth and depth

### 3.2.5. Temperature

The relationship between the impact depth and breaking volume and the standoff distance for various temperature is presented in Fig. 8. With the increase of temperature, the impact depth and the breaking volume increase. For hard and brittle rocks, uniaxial compressive strength and fracture toughness decrease obviously with the increase of temperature [58]. As the temperature of the rock increases, the impact of high-speed particles on the rock surface produces huge impact stress. Due to the reduction of the strength of the rock, the particle impact is more likely to produce tensile shear failure inside the rock, which is more conducive to the generation, expansion and extension of cracks inside the rock. Therefore, with the increase of the temperature, the impact depth of the rock and the volume of broken rock constantly increase.



**Fig. 8.** Temperature effect on the impact depth and depth

## 4. Conclusions

Establishment and confirmation of the breaking mechanism and optimal parameters of a particle jet for rocks

and other materials to be broken is extremely important. Therefore, the hot dry rock were impacted and broken by a particle jet under various parameters (such as velocity, diameter, concentration, confining pressure, and temperature) both by simulation. The results of this investigation are summarized as follows:

The rebound effect of particles constitutes an important factor in the crater diameter determination during a particle jet impact. The maximum value of the 1<sup>st</sup> maximum principal stress was 23.1 MPa, which was quite higher than the tensile strength of the rock (marble); therefore, the rock could easily be broken. The higher values of the 1<sup>st</sup> maximum principal stress were distributed on the surface areas of the crater during the impact progress.

The impact depth and breaking volume increased as the jet inlet velocity increased. The energy increase of particles was the main factor that influenced the rock breaking volume changes. Both the impact depth and breaking volume increased as the particle diameter increased. When the particle diameter exceeded 1.0 mm, the breaking volume began to increase rapidly. The breaking volume of a particle water jet strongly depended on the particle diameter.

Both the impact depth and breaking volume increased as the particle concentration increased. When the particle concentration was below a critical concentration (3%), the impact depth and breaking volume increased rapidly with the increase in the concentration. Moreover, when the particle concentration was  $>3\%$ , the impact depth and breaking volume increased slowly owing to the interference caused by the collisions between the incident and rebound particles. With the increase of confining pressure, the impact depth and breaking volume decrease. While with the increase of temperature, the impact depth and the breaking volume increase.

## Notations

### Nomenclature

$a$	volume correction	-
$A, B, N, C$	material strength parameters for HJC model	-
$C_b$	normalization constant	-
$C_e$	mixing efficiency coefficient	-
$C_p$	particle volume concentrations	%
$C_{MG}$	$u_s - u_p$ curve intercept	m/s
$D$	space dimensionality	-
$d_n$	nozzle diameter	cm
$D$	damage ratio	-
$D_i$	impact depth	mm
$E$	internal energy per unit volume	-
$E_p$	elasticity modulus of particles	-
$h$	SPH particle smooth length	-
$K_{jet}$	waterjet particle kinetic energy	J
$m_g$	gas mass flow rate	Kg/s
$m_p$	particle mass flow rate	Kg/s
$m_w$	water mass flow rate	Kg/s
$P_c$	confining pressure	MPa
$P_w$	nozzle pressure	MPa
$P_{MG}$	pressure	MPa
$P^*$	dimensionless hydrostatic pressure	-
$Q_w$	water flow rate	L/s
$S_1, S_2, \text{ and } S_3$	$u_s - u_p$ curve slope coefficient	-

$v_p$	particle inlet velocity	m/s
$v_w$	jet inlet velocity	m/s
$V_b$	rock breaking volume	cm <sup>3</sup>
$T_i$	temperature	°C
Greek letters		
$\gamma_0$	MIE-GRUNEISEN constant	-
$\dot{\epsilon}^*$	strain rate with a dimension of 1	-
$\rho_0$	Density	g/cm <sup>3</sup>
$\rho_p$	particle density	g/cm <sup>3</sup>
$\rho_w$	water density	g/cm <sup>3</sup>
$\sigma^*$	equivalent stress pressure	-
$\nu_p$	Poisson's ratio of particles	
$\Omega$	support domain of the point x	-

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