

## A STEAM/WATER COUNTERFLOW MODEL OF A VAPOUR DOMINATED GEOTHERMAL RESERVOIR WITH SURFACE RECHARGE & DEEP DRAINAGE

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### ABSTRACT

A 2 dimensional steam/water counterflow model of a Vapour Dominated Geothermal Reservoir has been developed. If liquid recharge is applied at the surface & deep drainage is available, a range of conditions can be established where natural chokes are formed in the reservoir. These chokes lead to a liquid dominated region overlying a vapour dominated region. The pressure (or depth) at which a choke will form is controlled by a narrow range of surface conditions. A selection of vertical sections have been constructed which model a series of possible reservoir conditions. These range from natural surface activity where the vapour dominated region is close to the surface through to the boundary of the vapour reservoir at its intersection with deep liquid dominated conditions. An interesting feature of the model is that vapour dominated conditions can first form at a pressure of 33 Bars(abs). This phenomenon appears to be unrelated to the maximum enthalpy of steam condition which occurs at a similar pressure.

### INTRODUCTION

Vapour Dominated Geothermal Reservoirs have a number of very interesting features, explanations for which, to date, do not appear to have been universally accepted. These features include:

(i) What appears to be a stable layer of water lying on top of a body of steam.

(ii) Steam reservoirs which always seem to occur at pressures in the region of 30 - 35 Bars(abs), with a pressure gradient close to steam static which in some reservoirs is continuous to indeterminate depth.

(iii) Production histories which can only be explained if relatively large amounts of Liquid water are present in the reservoir.

Explanations for the existence of a stable liquid condensate layer overlying a steam zone have been suggested by, among others, Schubert & Straus(1979 & 1980) and Schubert, Straus & Grant(1981).

The explanations referred to cite the Clapeyron equilibrium relationship coupled with relatively low rock permeabilities as providing a suitable mechanism.

James(1968) and McNitt(1977) have described the formation of natural steam zones in terms of the minimum enthalpy of saturated steam as both features seem to occur at similar pressures.

Mounfort(1980) has estimated the water content of the Kamojang reservoir in terms of residual saturations with a small quantity of mobile water present, and Grant(1979) has carried out a correlation based on the transient response of gaseous contaminants in the steam.

The classic model of a vapour dominated reservoir is that proposed by White, Muffler & Truesdell(1971) which relies on an overlying cap rock

The author has developed a reservoir model where, if recharge water is added at the surface and a counterflow conditions are assumed with steam rising from depth, a range of conditions occurs where a natural choke will be established which effectively restricts the upward flow of steam and the downward flow of water. Liquid dominated conditions can thus occur above the choke, and vapour dominated conditions below. The formation of this choke, which appears to be independent of geological features or rock permeabilities leads to reservoir models that are remarkably similar to real reservoirs.

Features of the model include-

(i) Formation of natural boundaries which can be in a uniformly permeable conditions purely is a function of the properties of steam & water.

(ii) Natural surface activity can occur in the same uniform conditions without affecting the other features of the model.

(iii) The first formation of vapour dominated conditions occurs at a pressure of 33 Bars(abs), and this feature seems to be independent of the maximum enthalpy of steam condition.

(iv) An estimate of actual reservoir permeability leads directly to the calculation of recharge and hence the steam/water proportions throughout the reservoir, and the full range of reservoir simulation calculations can be performed.

A conceptual model is first developed where the mass & energy balance required to maintain thermodynamic equilibrium throughout a vertical section is calculated. Determination of a range of possible conditions leads to a two dimensional model that is modified to approximate real conditions. Finally, the conceptual model is corrected so that gravity, viscosity and relative permeability effects are included.

### CONCEPTUAL MODEL

If one considers the standard model of a vapour dominated geothermal reservoir, overlain by a condensate layer and underlain by a deep brine level, the vertical section is shown in Fig 1.

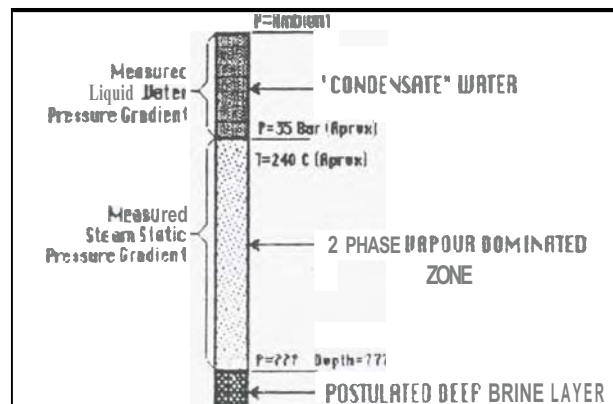


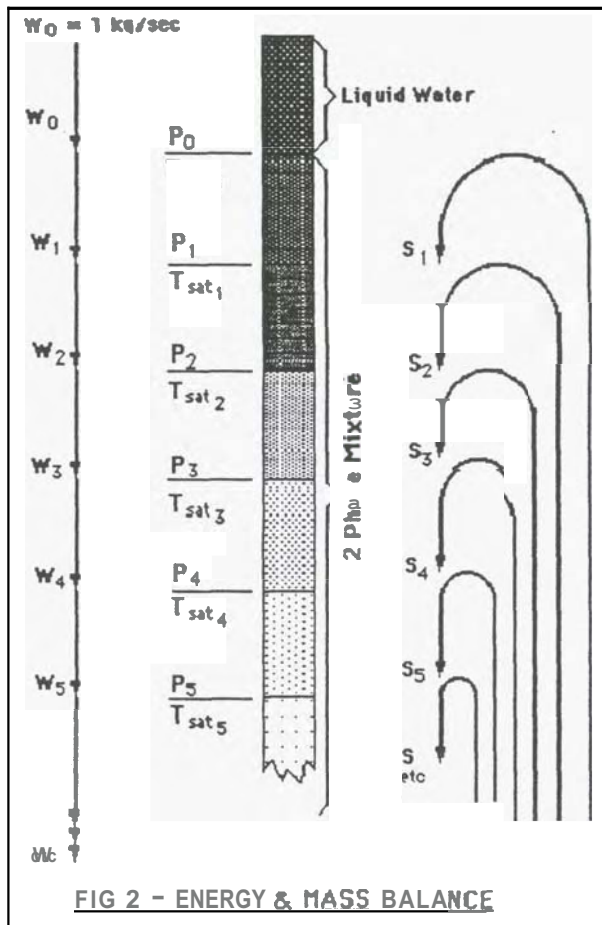
FIG 1-VAPOUR DOMINATED RESERVOIR-VERTICAL SECTION

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If such a system is in a state of thermodynamic equilibrium throughout its depth, and a counterflow model is assumed then the steam/water proportions and relative flow rates can be calculated for any point on the vertical section.

Consider a unit inflow of water at the top of the system. This water will be heated by conduction & convection until it attains the saturation temperature for depth. Upon reaching this temperature, the water can co-exist with steam rising from below. As the water continues to descend through the rising steam it must gradually increase in temperature so that saturation conditions are maintained. Because the water is in intimate contact with the steam, and all temperatures, including the rock temperature, are controlled by boiling point for depth conditions, the water will be heated (with a corresponding increase in enthalpy) either directly or indirectly, by the condensation of steam.

The separate steam & water masses found at any depth (equivalent to pressure) in the reservoir, based on a unit inflow of water at the top of the system can be calculated by performing a sequential mass & balance energy balance for every pressure interval as shown in Fig 2



Let  $P_0$  be the pressure where the unit inflow of water,  $W_0$ , first attains saturation temperature,  $T_{sat0}$ .

Then the mass of steam,  $S_1$ , condensed as  $W_0$  passes through the pressure interval  $P_0 - P_1$  and increases in temperature from  $T_{sat0}$  to  $T_{sat1}$  is approximately

$$S_1 = W_0(h_{f1} - h_{f0})/h_{fg1}$$

$$\& \quad W_1 = W_0 + S_1$$

and similarly,

$$S_2 = W_1(h_{f2} - h_{f1})/h_{fg2} + S_1 - S_1(h_{g1}/h_{g2})$$

$$\& \quad W_2 = W_1 + S_2 + S_1 - S_1(h_{g1}/h_{g2})$$

where

$$S_1 - S_1(h_{g1}/h_{g2})$$

is a correction made to the steam & water flows to account for the fact that water will be condensed out or evaporated by the rising steam if the pressure is below or above that where the maximum enthalpy of steam occurs (31.19 Bar)

Steam & water masses, at every pressure interval, can then be calculated in a similar way as follows-

$$S_3 = W_2(h_{f3} - h_{f2})/h_{fg3} + (S_1 + S_2)(1 - (h_{g2}/h_{g3}))$$

$$W_3 = W_2 + S_3 + (S_1 + S_2)(1 - (h_{g2}/h_{g3}))$$

$$S_4 = W_3(h_{f4} - h_{f3})/h_{fg4} + (S_1 + S_2 + S_3)(1 - (h_{g3}/h_{g4}))$$

$$W_4 = W_3 + S_4 + (S_1 + S_2 + S_3)(1 - (h_{g3}/h_{g4}))$$

$$S_5 = W_4(h_{f5} - h_{f4})/h_{fg5} + (S_1 + S_2 + S_3 + S_4)(1 - (h_{g4}/h_{g5}))$$

$$W_5 = W_4 + S_5 + (S_1 + S_2 + S_3 + S_4)(1 - (h_{g4}/h_{g5}))$$

etc, down to the critical point, if necessary, and for a range of values of  $P_0$ .

A typical table of results is as shown in part, in Fig 3, where, in this example,  $P_0 = 3$  Bars.

P. (Bar)	Steam Condensed (kg/s)	Water Loss/Gain (kg/s)	Water Flow (kg/s)	Steam Flow (kg/s)
1.0000			1.0000	0.0000
1.2000			1.0000	0.0000
1.4000			1.0000	0.0000
1.6000			1.0000	0.0000
1.8000			1.0000	0.0000
2.0000			1.0000	0.0000
2.2000			1.0000	0.0000
2.4000			1.0000	0.0000
2.6000			1.0000	0.0000
2.8000			1.0000	0.0000
3.0000			1.0000	0.0000
3.2000	0.0044	0.0000	1.0044	0.0044
3.4000	0.0042	0.0000	1.0086	0.0086
3.6000	0.0040	0.0000	1.0126	0.0126
3.8000	0.0039	0.0000	1.0166	0.0166
4.0000	0.0038	0.0000	1.0203	0.0203
4.2000	0.0036	0.0000	1.0239	0.0239
4.4000	0.0035	0.0000	1.0274	0.0274
4.6000	0.0034	0.0000	1.0309	0.0309
4.8000	0.0033	0.0000	1.0342	0.0342
5.0000	0.0032	0.0000	1.0374	0.0374
5.2000	0.0032	0.0000	1.0405	0.0405
5.4000	0.0031	0.0000	1.0436	0.0436
5.6000	0.0030	0.0000	1.0466	0.0466
5.8000	0.0030	0.0000	1.0495	0.0495
6.0000	0.0029	0.0000	1.0524	0.0524
6.2000	0.0028	0.0000	1.0552	0.0552
6.4000	0.0028	0.0000	1.0579	0.0579
6.6000	0.0027	0.0000	1.0606	0.0606
6.8000	0.0027	0.0000	1.0632	0.0632
7.0000	0.0026	0.0000	1.0658	0.0658
8.0000	0.0124	-0.0001	1.0781	0.0781
9.0000	0.0115	-0.0001	1.0895	0.0895
10.0000	0.0108	-0.0001	1.1002	0.1002
11.0000	0.0102	-0.0001	1.1103	0.1103
12.0000	0.0097	-0.0001	1.1198	0.1198
13.0000	0.0093	-0.0001	1.1290	0.1290

FIG 3  
STEAM & WATER MASSES

$P_0 = 3$  BAR

if the volume occupied by the steam/water mixture at each pressure interval is calculated ie,

$$V_n = W_n \cdot v_{fn} + (S_1 + S_2 + S_3 + \dots + S_n) \cdot v_{gn}$$

and if this process is repeated over all pressure intervals and for a range of values of  $P_0$ , then a set of results are obtained which can be plotted as the family of curves shown in Fig 4, where each curve represents a particular value of  $P_0$



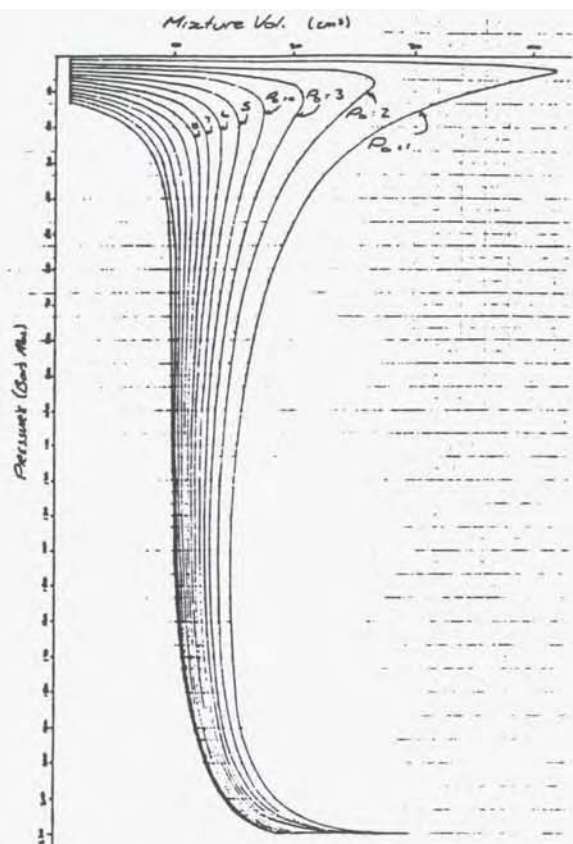


FIG 4

### STEAM/WATER MIXTURE VOLUME vs PRESSURE FOR A RANGE OF VALUES OF $P_0$

An interesting phenomenon can be noted from Fig 4. For values of  $P_0$  between 1 Bar(abs) & 12 Bar(abs), the volume required to contain the equilibrium steam & water flows starts at a small value (equal to the specific volume of water) at  $P_0$ , rises to a maximum value at some higher pressure, and then falls off. For values of  $P_0$  greater than 12 Bar(abs), a similar maximum does not occur and the volume required continues to increase steadily with increasing pressure.

For values of  $P_0$  between about 2 Bars(abs) & 12 Bars(abs) the volume increases again at higher pressures and eventually becomes larger than the previously determined maximum value.

The occurrence of the upper maximum, which is controlled by a relatively small range of possible values of  $P_0$  is suggested to be the feature which leads to the formation of vapour dominated reservoirs and steam caps above exploited liquid dominated reservoirs.

If it is assumed that the reservoir has uniform permeability (any value), and the calculated value for the maximum volume is divided by unit thickness, then an area is found which is just sufficient to pass the upward steam and the downward water flows required to keep the reservoir in equilibrium. (Based on a recharge at the top of the system of 1 kg/s). As the reservoir has uniform permeability, then this area is also available above and below the point where the maximum occurs. At all points where the total available area is greater than that required to pass the equilibrium flow rates, the excess area can fill with either steam or water.

At points in the reservoir above where the maximum required area occurs, the supply of steam from below is choked by the controlling reservoir area and so the excess area must fill with water.

Conversely, below the maximum, the supply of water is choked and so the excess reservoir area, or volume, will fill with steam and a Vapour Dominated Reservoir will result. This vapour dominated condition will continue downwards until the reservoir area required to pass the equilibrium flow rates becomes greater

than that available at the upper maximum. (Although further development of the model will show that the base water level is probably controlled by the boundary condition where the upper maximum and the lower increase above this controlling value become coincident).

The above argument applies to reservoirs of uniform permeability. We can briefly consider the effects that would occur in a real reservoir where the permeability is far from uniform. What happens in this case is that the controlling permeability is the *minimum* permeability (which can still have any value), and all other permeabilities found in the reservoir will be greater than this minimum. Where greater permeabilities occur above the choke point, then even more water can accumulate, and if greater permeability than the minimum occurs below the choke point then more steam can accumulate.

For values of  $P_0$  greater than about 12 Bars(abs) where no maximum occurs, the reservoir area required to pass the equilibrium flow rates is always less than the controlling area (where ever it may occur at depth) so the excess area available will always fill with water and so Vapour Dominated conditions will not be encountered anywhere in the reservoir.

Referring back to Fig 4, for each value of  $P_0$  considered, the maximum volume (or area) can be determined and the difference between this volume and that required at every pressure interval above and below calculated and the excess volume filled with water or steam as the case may be.

The density of these new mixtures (containing either more water or more steam than the equilibrium amounts previously calculated) can then be determined and hence the depth of every pressure interval calculated with respect to the top of the system. The results are plotted in Fig 5.

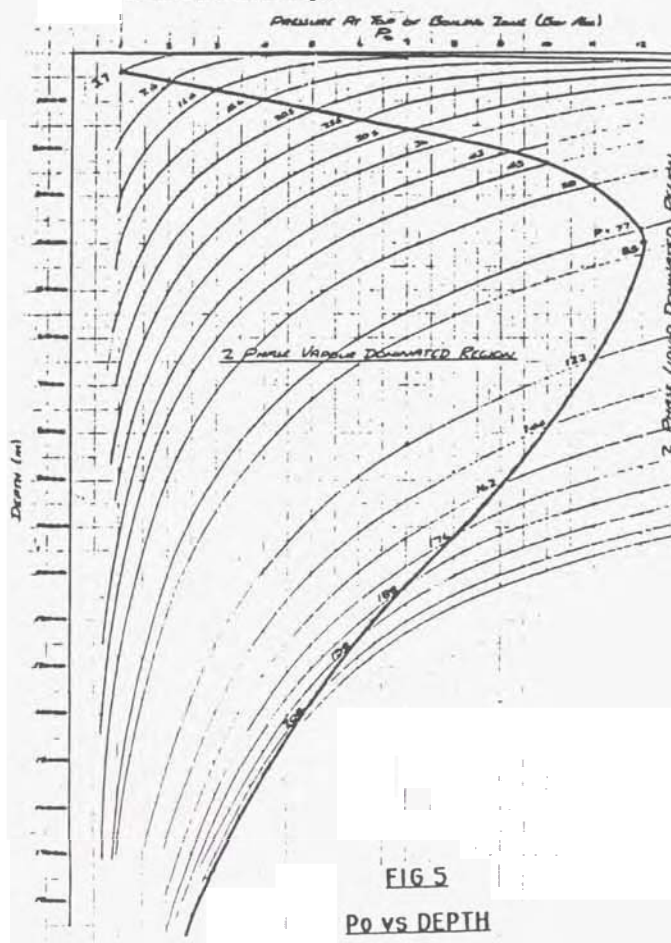


FIG 5

### $P_0$ vs DEPTH

Points to note from Fig 5 (apart from the fact that it bears no resemblance to my known vapour dominated reservoir) are as follows-



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(i) The heavy line is the locus of the depths where the upper maximum area or the lower area equal to this upper maximum occurs. The upper & lower maximums become coincident when  $P_0$  is just over 12 Bars(abs) and at a pressure of about 85 Bars(abs).

(ii) The lighter lines are isobars drawn through the pressure corresponding to the upper maximum area and the lower area corresponding to this for each value of  $P_0$ .

(iii) Note that the terms *liquid dominated* and *vapour dominated* differ slightly from the normal convention. The *Liquid Dominated Region* is where excess reservoir volume will fill with water. The *Vapour Dominated Region* is where excess reservoir volume will fill with steam. Both of the above regions contain a 2 phase mixture of steam & water. A single phase liquid water zone occurs as a thin wedge at the top of the system.

There are two modifications that can be made to the reservoir cross section shown in Fig 4 in order to make it more realistic. Firstly, consider the boundary at  $P_0=12$  Bars(abs). This is the shallowest depth at which an external influence can affect the vapour region, as at every other point to the left on the diagram, steam & water flows and their relative proportions are controlled by the choke that occurs at the upper area maximum. Thus, the pressure gradient within the vapour dominated region cannot be affected by what is happening outside this region (within the range of conditions that allow a vapour region to form in the first place). However, when  $P_0=12$  Bars(abs), the lower and upper maximum area locations become coincident. Now, at the lower area restriction, similar conditions to those applying at the upper maximum occur, i.e. there is just sufficient reservoir area available to pass the equilibrium steam & water flows. Below this point there is insufficient area to either the steam or the water flow will be restricted. If the water flow is restricted then the reservoir will gradually fill up with water and eventually the vapour dominated reservoir will disappear. Therefore, the steam flow must be restricted and the water allowed to drain. Thus the lower reservoir area restriction will roughly approximate the deep water level with most of the steam being produced by boiling from this surface.

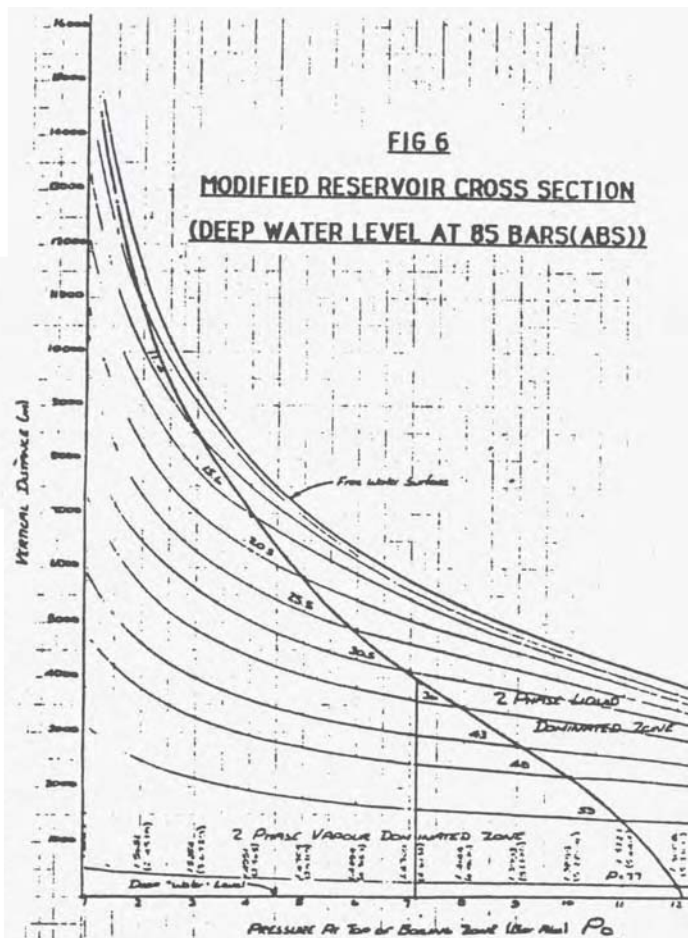
From Fig 5, this deep water level will occur at a pressure of approximately 35 Bars(abs). Thus, the first modification we can make to Fig 5 is to reconstruct the 35 Bars(abs) isobar and make it horizontal. The result is shown in Fig 6, which is a slightly more realistic looking reservoir cross section.

The second modification that can be made is to ignore regions of the vapour dominated region with pressures less than 31.19 Bars(abs) (corresponding to the maximum enthalpy of steam), simply because vapour dominated reservoirs do not seem to occur at pressures below this value. This point occurs at approximately  $P_0=7.2$  Bars(abs). The reservoir is then made symmetrical about this value of  $P_0$  and the cross section as shown in Fig 7 plotted.

The reservoir cross section plotted in Fig 7 has a number of encouraging features. Among these are

(i) The existence of a vapour dominated reservoir is governed only by the thickness of the relatively thin top liquid water layer. If the thickness of this layer is outside certain limits it appears the vapour dominated region collapses into two phase liquid dominated conditions. In other words, there seems to be no dependence on geological conditions or rock permeabilities. (Although these features can play a significant role in modifying the standard model)

(ii) The model has naturally occurring boundaries, created by the properties of steam alone



(iii) The conditions that determine the maximum depth of the deep water layer have been established. (As the existence of the vapour dominated region depends on there being adequate drainage for the recharge water & condensate, a higher zone of impermeable rock will lead to a reduction in this maximum possible depth)

(iv) The slope of the isobars on top of and to the sides of the vapour dominated region will enable excess recharge water to flow away, which is important as the amount able to enter the vapour dominated region is strictly controlled by the maximum area limitation.

(v) The heat source is centrally located and the slope of the isobars in the vapour dominated region will allow the rising steam some horizontal movement towards the boundaries.

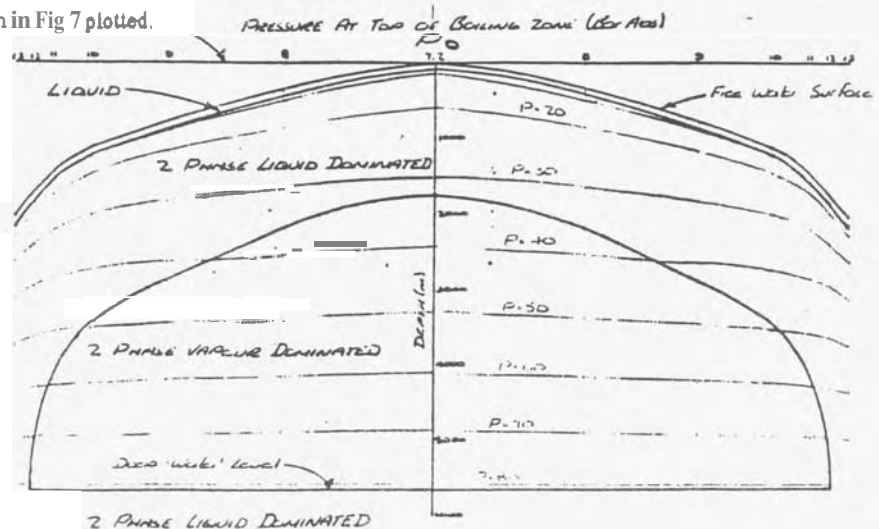


FIG 7

DEVELOPED CROSS SECTION OF RESERVOIR





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$$(w_s v_s - w_w v_w) K_{rw}^2 + 2. (w_w v_w) K_{rw} - (w_w v_w) = 0$$

Which has one root of  $K_{rw}$  between zero & one. Substituting this value in 6 gives the value of  $KA_{min}$  at each pressure interval.

This calculation can be repeated for all pressure increments and over a range of values of  $P_0$  and a family of curves of  $KA_{min}$  vs  $P$  similar to those shown in Fig 4 obtained.

Curves plotted for low values of  $P_0$  have the same characteristic maximum value as found in the conceptual model (in this use  $KA$ ) which can be used to recalculate  $K_{rw}$  at all pressure increments above and below the interval where this maximum value of  $KA_{min}$  occurs.

#### Rearranging 6

$$KA(P_w - P_0) K_{rw}^2 + (w_s v_s - w_w v_w - KA(P_w - P_0)) K_{rw} + w_w v_w = 0$$

Which when solved gives the two values of  $K_{rw}$  which occur when an  $nm$  greater than the minimum required to pass the equilibrium flows is available. Thus, at all pressure intervals above and below the interval where  $KA$  is a maximum, there is a choice of values of  $K_{rw}$  - representing either a higher or lower volume fraction of water than occurs when  $KA = KA_{min}$ . Using the same logic that was used in the argument for the conceptual model, the higher value of  $K_{rw}$  (ie higher water content) can be chosen for intervals above where  $KA$  is a maximum, and the lower value of  $K_{rw}$  (ie higher steam content) below.

The revised reservoir model can thus be developed using exactly the same procedures as that developed for the conceptual model.

Vapour dominated conditions occur over only a very small range of values of  $P_0$  is between  $P_0 = 1$  &  $P_0 = 2.6$  Bars(abs) approximately but the reservoir model produced is extremely interesting.

Calculations as detailed as those performed when constructing the conceptual model have not yet been completed. The model produced therefore lacks the resolution, but sufficient information is available to plot the reservoir section shown in Fig 10. Fig 10 is equivalent to Fig 6 of the conceptual model but includes additional detail below the horizontal isobar down to a pressure of 60 Bars(abs).

In Fig 10, the upper & lower  $KA$  maximums are coincident at a pressure of 33 Bars(abs) approx, thus, this isobar has been drawn horizontally as in the conceptual model. Similar interpretations as made using the conceptual model can be made to this model, but with the following important modifications or exceptions;

(i) The occurrence of the coincident upper & lower  $KA$  maximums at 33 Bars(abs) is extremely significant as this is where the first formation of vapour dominated conditions will take place in the reservoir. The formation of natural steam which occurs at about this pressure in real reservoirs therefore may not be due to some property related to the maximum enthalpy of steam (which occurs at a pressure of 31.19 Bars(abs)), but could arise naturally as the result of counter flow reservoir conditions.

(ii) We do not need to discard half the model, as had to be done with the conceptual model, but instead can allow some surface activity in the form of surface water close to the boiling point of steam vents which can occur at  $P_0 = 1$  or less.

(iii) The natural boundaries still occur but in order for this model to fit a real reservoir, the theory put forward in the conceptual model concerning the position of the deep water level has to be modified somewhat in order that the vapour dominated region can have a reasonable thickness at pressures at or above 33 Bars(abs). It will be noted from Fig 10, that when the 33 Bar isobar is made

tend to be drained away from the vapour dominated region ie conditions outside the reservoir can have no effect. However, below the horizontal isobar, the isobar slope will allow external pressure conditions to affect the reservoir and the reservoir pressure will reflect conditions outside ie liquid dominated. This effect will occur for any isobar that is made horizontal. We could, in theory, have a horizontal water surface at the critical point, 221 Bars(abs). In this case there would be a very large vapour dominated reservoir with a higher pressure gradient than its surroundings. In practice though, it seems that the deep water level will occur just above the deepest permeable drainage zone. Further investigation of the deep reservoir may reveal a maximum possible depth, but unfortunately, it is not possible at the moment to make a prediction based on the properties of a uniformly permeable reservoir.

(iv) An actual value of  $KA$  is calculated for each value of  $P_0$  considered. Substitution of a value for average rock permeability will allow the reservoir  $nm$  required to accept a recharge of 1 kg/sec to be readily estimated. This figure can be used in a range of reservoir engineering calculations.

(v) Fig 11 is an extract of the data produced for  $P_0 = 2.0$  Bars(abs). It will be noted that in the Vapour Dominated region (between  $P = 9$  & 60 Bars(abs)), the value of  $K_{rw}$  (the water volume fraction in effect) lies approximately in the range .34-.38. Thus the pressure gradient in this region is about 30% of the equivalent hydrostatic pressure gradient.

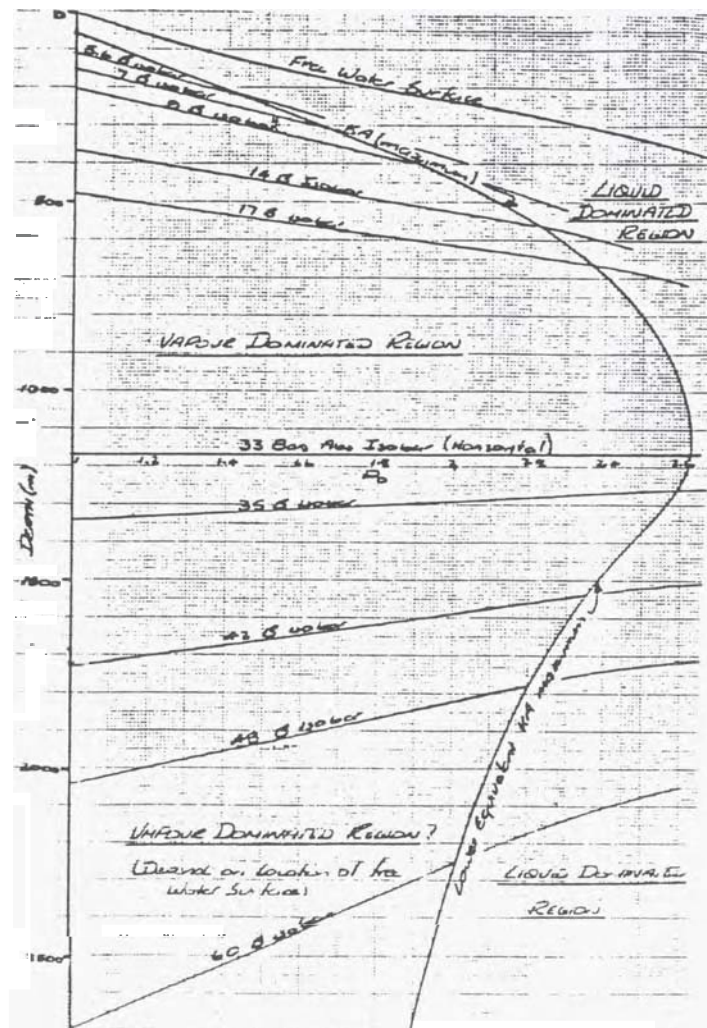


FIG 10

CROSS SECTION OF A POSSIBLE REAL RESERVOIR

P (Bar)	K <sub>A</sub> (m <sup>4</sup> )	K <sub>rw</sub> (Liq. Dom.)	K <sub>rw</sub> (Vap. Dom.)	ρ (Pa/m)	Depth (m)
1	1.274E-10	1.0000	0.2374	7143.1029	3.1253
1.2		1.0000	0.2273	7212.1778	6.2743
1.4		1.0000	0.2210	7247.2403	9.4616
1.6		1.0000	0.2144	7287.3884	12.6645
1.8		1.0000	0.2082	7324.8267	15.9437
2		1.0000	0.2208	7189.5792	19.3182
2.2	5.967E-11	0.9316	0.2312	7075.8883	22.8016
2.4	7.365E-11	0.8745	0.2409	6970.5146	26.3933
2.6	8.325E-11	0.8275	0.2494	6877.3052	30.0899
2.8	9.045E-11	0.7875	0.2601	6765.6280	33.9049
3	9.68E-11	0.7502	0.2683	6677.0784	37.8284
3.2	1.018E-10	0.7187	0.2763	6591.8958	41.8614
3.4	1.053E-10	0.6928	0.2804	6542.3502	45.9841
3.6	1.078E-10	0.6738	0.2906	6438.5091	50.2335
3.8	1.111E-10	0.6478	0.2979	6360.9817	54.5956
4	1.136E-10	0.6263	0.3061	6276.4767	59.0782
4.2	1.158E-10	0.6069	0.3125	6208.0049	63.6726
4.4	1.174E-10	0.5915	0.3149	6176.6983	68.3529
4.6	1.183E-10	0.5812	0.3172	6146.8932	73.1190
4.8	1.189E-10	0.5739	0.3244	6073.1057	78.0068
5	1.204E-10	0.5581	0.3336	5982.3628	83.0335
5.2	1.218E-10	0.5425	0.3391	5924.5169	88.1746
5.4	1.227E-10	0.5305	0.3459	5855.8404	93.4422
5.6	1.235E-10	0.5197	0.3478	5830.8133	98.7986
5.8	1.238E-10	0.5137	0.3509	5794.8734	104.2555
6	1.242E-10	0.5085	0.3481	5811.8672	109.7629
6.2	1.24E-10	0.5092	0.3582	5715.7376	115.4506
6.4	1.25E-10	0.4943	0.3654	5644.4580	121.2386
6.6	1.256E-10	0.4837	0.3739	5562.7699	127.2015
6.8	1.262E-10	0.4721	0.3816	5488.0085	133.3162
7	1.266E-10	0.4617	0.3722	5538.5292	139.4691
8	1.263E-10	0.4660	0.4151	5135.6998	177.0348
9	1.274E-10	0.4151	0.3774	4781.7407	201.9981
10	1.267E-10	0.4531	0.3820	4843.5216	227.0428
11	1.269E-10	0.4434	0.3839	4811.4804	252.6569
12	1.269E-10	0.4447	0.3777	4726.3761	279.1420
13	1.267E-10	0.4524	0.3803	4733.9317	305.9939
14	1.268E-10	0.4493	0.3697	4608.8603	333.9946
15	1.263E-10	0.4620	0.3614	4507.8793	363.0524
16	1.258E-10	0.4720	0.3673	4536.3289	392.3548
17	1.261E-10	0.4667	0.3596	4446.7519	422.6831
18	1.256E-10	0.4757	0.3540	4352.9103	454.1102
19	1.252E-10	0.4853	0.3546	4333.3718	486.1259
20	1.252E-10	0.4860	0.3546	4303.5517	518.8135
21	1.251E-10	0.4880	0.3512	4240.3611	552.4450
22	1.247E-10	0.4942	0.3432	4122.1283	587.5109
23	1.239E-10	0.5073	0.3459	4131.8330	622.9633
24	1.241E-10	0.5047	0.3458	4108.6975	659.0866
25	1.241E-10	0.5062	0.3471	4104.8673	695.7155
26	1.241E-10	0.5053	0.3436	4044.6178	733.3689
27	1.237E-10	0.5115	0.3446	4009.3239	771.8369
28	1.236E-10	0.5147	0.3451	3999.5221	810.8833
29	1.237E-10	0.5147	0.3501	4012.7746	850.2835
30	1.24E-10	0.5118	0.3485	3979.1901	890.5029
31	1.238E-10	0.5149	0.3506	3960.3292	931.4029
32	1.239E-10	0.5161	0.3467	3901.5857	973.4152
33	1.234E-10	0.5225	0.3506	3904.0419	1015.8971
34	1.236E-10	0.5210	0.3509	3895.8415	1058.9656
35	1.237E-10	0.5210	0.3559	3911.2493	1102.3596
36	1.24E-10	0.5179	0.3548	3888.4056	1146.5066
37	1.239E-10	0.5198	0.3589	3894.9106	1191.0772
38	1.241E-10	0.5180	0.3589	3870.0440	1236.4346
39	1.24E-10	0.5201	0.3583	3838.2812	1282.6719
40	1.239E-10	0.5233	0.3664	3868.9623	1335.5906
42	1.245E-10	0.5173	0.3724	3872.5739	1429.0418
44	1.248E-10	0.5148	0.3764	3858.2612	1503.9699
46	1.249E-10	0.5147	0.3841	3878.8367	1579.4993
48	1.254E-10	0.5101	0.3957	3935.5468	1654.9246
50	1.26E-10	0.5005	0.4071	3984.9059	1730.3878
52	1.265E-10	0.4918	0.4168	4017.5541	1806.2019
54	1.268E-10	0.4854	0.4377	4135.1036	1880.7975
56	1.273E-10	0.4671	0.4280	4061.3063	1957.7024
58	1.271E-10	0.4753			
60	1.277E-10				

FIG 11

## RESULTS TABLE FOR Po= 2.0 Bars(abs)

Space does not permit a discussion concerning the apparent disappearance of the condensate layer. however, the author is currently re-evaluating measurements taken in the so called condensate & vapour regions of some real reservoirs.

The difficulty of interpreting downwell pressure & temperature measurements taken in two phase reservoirs is well known. Interpreting these measurements by comparing them with the theoretical well conditions described in this paper as opposed to trying to describe the reservoir in terms of the measurements is leading to some very interesting observations which will be reported in a future paper

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## NOTATION

A	Area (m <sup>2</sup> )
g	Acceleration due to gravity (~9.81 m/s <sup>2</sup> )
h <sub>l</sub>	Specific enthalpy of liquid (kJ/kg)
h <sub>g</sub>	Specific enthalpy of vapour (kJ/kg)
h <sub>fg</sub>	Latent heat of evaporation (kJ/kg)
K	Permeability (m <sup>2</sup> )
k <sub>rw</sub>	Relative permeability with respect to water
k <sub>rs</sub>	Relative permeability with respect to steam
P	Pressure (Bar or Pa)
P <sub>0</sub>	Pressure at first water saturation temperature (Bar)
S <sub>n</sub>	Steam condensed in each pressure interval (kg/s)
T	Temperature (C)
U <sub>w</sub>	Mass flux density of water (kg/m <sup>2</sup> s)
U <sub>s</sub>	Mass flux density of steam (kg/m <sup>2</sup> s)
v <sub>f</sub>	Specific volume of water (m <sup>3</sup> /kg)
v <sub>g</sub>	Specific volume of steam (m <sup>3</sup> /kg)
V	volume (m <sup>3</sup> )
W <sub>n</sub>	Water flow (kg/s)
W <sub>w</sub>	Water flow (kg/s)
W <sub>s</sub>	Steam flow (kg/s)
Δ	Change in
ν <sub>w</sub>	Kinematic viscosity of water (m <sup>2</sup> /s)
ν <sub>s</sub>	Kinematic viscosity of steam (m <sup>2</sup> /s)
ρ <sub>w</sub>	Density of water (kg/m <sup>3</sup> )
ρ <sub>s</sub>	Density of steam (kg/m <sup>3</sup> )

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